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Advancements in Nanodevice Design for Label-Free Detection of Biomolecules

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Abstract:

In the realm of nanotechnology, the detection of biological molecules poses both challenges and opportunities. This paper introduces a novel device structure, the CPB DM-JLTFET, designed to address complications encountered during nanodevice manufacturing, such as temperature coefficient variations and dopant concentration fluctuations. Leveraging charge plasma principles, this innovative structure incorporates an oxide substance in the etching process, resulting in an underlap zone within the nanodevice. Through label-free detection techniques reliant on the electrical characteristics of biomolecules, the CPB DM-JLTFET demonstrates superior performance attributes, including enhanced sensitivity and reduced leakage current compared to existing FET devices. A comprehensive analysis of device factors, including cavity length and thickness, reveals their significant influence on detection sensitivity. Utilizing ATLAS software for simulation, this research underscores the potential of the CPB DM-JLTFET in advancing the development of cost-efficient biomedical diagnostic tools.

Keywords: Biomolecules, TFET, transistor, semiconductor, nano

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1. INTRODUCTION

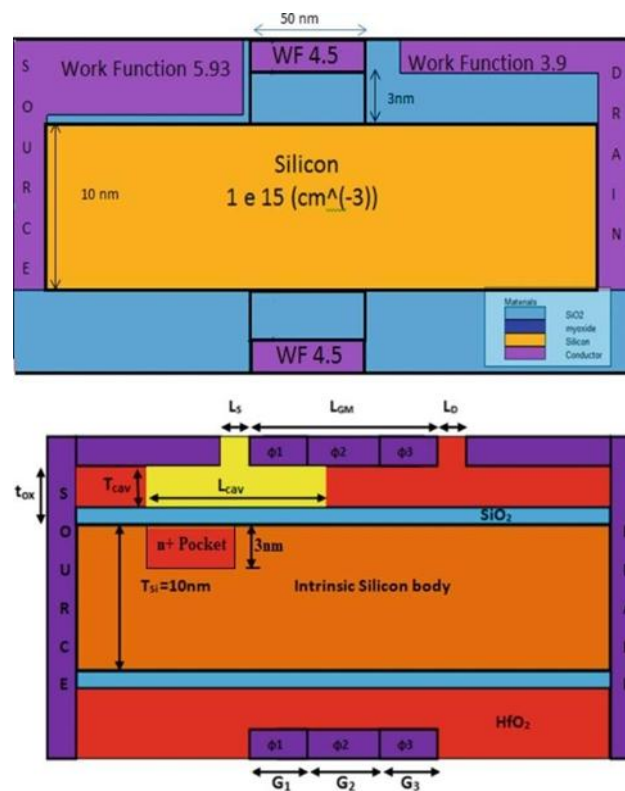
Biosensors (Naresh & Lee, 2021), analytical instruments capable of detecting target molecules without invoking a reaction system, hold immense potential for high-sensitivity biomolecule detection (Malhotra et al., 2016). Contrary to conventional wisdom, the significance of biosensors amplifies with decreasing biomolecule size, especially in the realm of high-target biosensors (HTB) (Iwanaga, 2021) utilizing dielectric and charge density principles (Stobiecka et al., 2017). This study focuses on harnessing dielectric constants within a Field-Effect Transistor (FET) (Reddy & Panda, 2021) biosensing framework to achieve precise biomolecule detection. However, challenges such as temperature sensitivity, real-time analysis, and Random Dopant Fluctuations (RDF) pose significant obstacles to practical biosensor implementation (Nguyen et al., 2017). To address these challenges, meticulous consideration of surface potential, spacer length, thickness, and cavity alignment becomes imperative in sensor design.

Proposed advancements in vertical device architecture aim to mitigate Short Channel Effects (SCE) (Zanoni et al., 2024) by incorporating multiple gates (Ajay et al., 2015). The immobilization of charged biomolecules beneath the gate nanocavity alters the dielectric constant at the semiconductor-oxide interface, manifesting discernible changes in electrical specifications. Consequently, understanding the influence of ordered parameter disparities on sensitivity and device current becomes paramount. Traditional TFET fabrication processes, reliant on high process-ion and high-temperature annealing thermal budgets, encounter formidable dopant fluctuation issues, exacerbating RDF-induced current fluctuations (Kanungo et al., 2016).

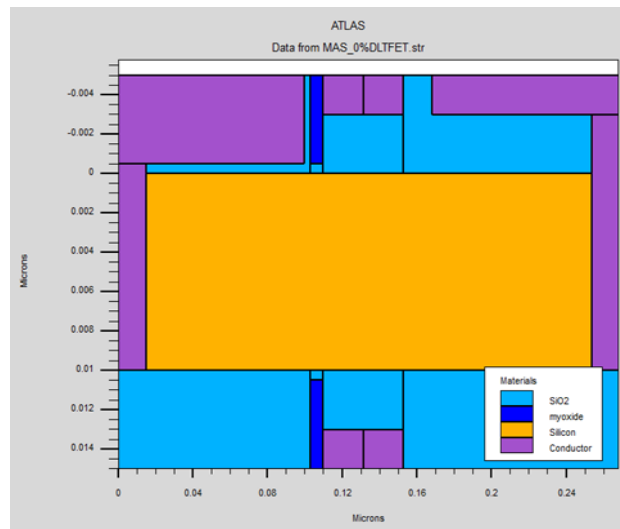
To surmount these obstacles, this study employs Plasma and electrostatic technologies, with Plasma doping introducing the charge plasma doping JLTFET. The integration of a dual-gate structure, comprising distinct gate electrodes with varied work functions, enhances device performance. The JLTFET design resolves manufacturing challenges while augmenting subthreshold slope. Notably, the JLTFET gate underlay facilitates the engraving of a hollow area, enriching the cavity region for biomolecule attachment.

Moreover, the study proposes the double-gate dielectric-modulated JLTFET (DGDMJLTFET) biosensor, featuring a nanocavity space above the tunneling junction for biomolecule immobilization. Wet etching techniques are employed to create the nanocavity, albeit concerns regarding biomolecule partial filling and steric resistance necessitate thorough simulation and analysis. By evaluating drain current, ION/IOFF ratio, and subthreshold sensitivity to both neutral and charged biomolecules across various dielectric constants, this study elucidates the DGDMJLTFET's superior sensitivity and performance in biomolecule detection.

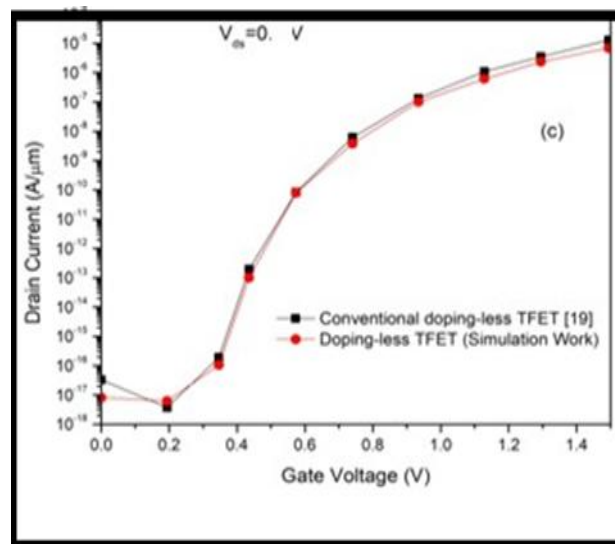
2. PROPOSED STRUCTURE



(a)



(b)



(c)

Figure 1. Structures (a) TFET (b) Underlap dielectric modulated JLTFET incorporating a charge-plasma-based gate, (c) Comparison.

Figure 1 depicts the features curve for TFETs (Pown & Lakshmi, 2020) alongside the innovative CPB DM-JLTFET structure. The design incorporates specific metal work function electrodes to induce source and drain regions within an ultrathin silicon film. A block of platinum is utilized to generate holes on the source side in conjunction with an intrinsic silicon semiconductor.

Imitation include the Debye length, determined by factors. Notably, the gate overlap area and gate underlap region are critical aspects of the channel. The intrinsic one, set at approximately $1.0 \times 10^{15} \text{ cm}^{-3}$, is employed for device characterization.

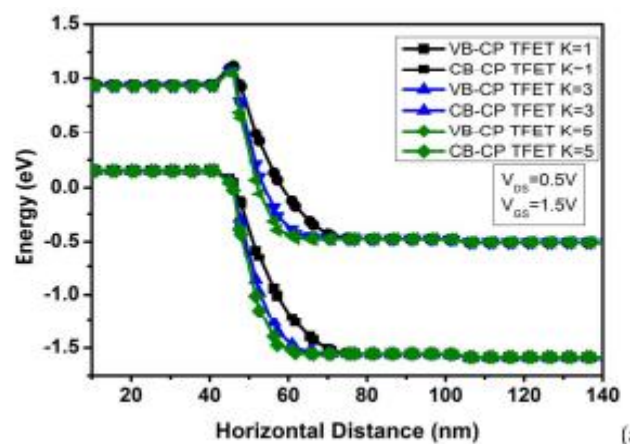
The opening area distance, ranging from 5 nano to 7 nm, corresponds to lengths of gates underlap. SiO₂ sheet separates electrodes, while drain-gate spacers measure 15 nm in thickness and source-gate spacers 3 nm. Results closely resemble those of the standard device, except for the additional gate underlap area.

3. RESULTS AND DISCUSSION

In modelling the electrical properties of the device, consideration is given to the dielectric constant of the gate underlaps and overlaps. Additionally, for charged biomolecules, the fixed charge of the SiO₂ interface area is accounted for.

3.1 Influence on energy bands

Fig. 2(a) and 3 depict the impact of particles. The formation of cavities within the transistor elevates the barrier, thereby reducing the likelihood of electron tunnelling and subsequently decreasing current conduction (Ghosh & Akram, 2013).



(a)

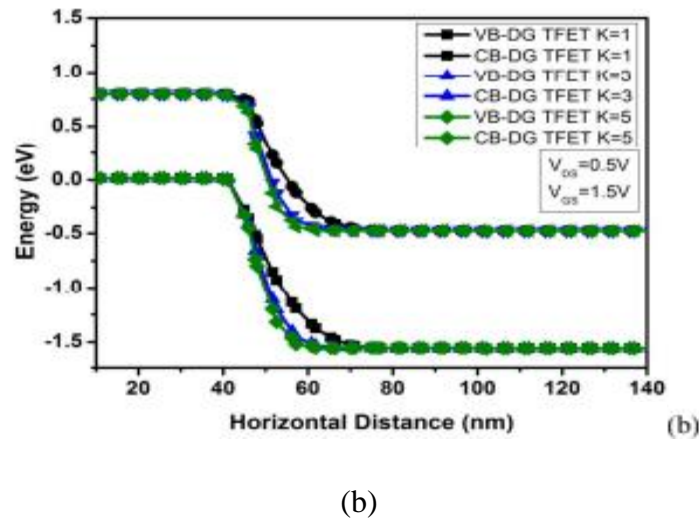


Figure 2. Alteration of conduction and valence bands due to (a) neutral biomolecules with values of K maintained at one, three and five (b) influence of living particles with charge ($N_f = -5e11$ and $5e11$) at $V_{gs} = 1.5V$; $V_{ds} = 0.5V$, affecting the 5 nm cavity length.

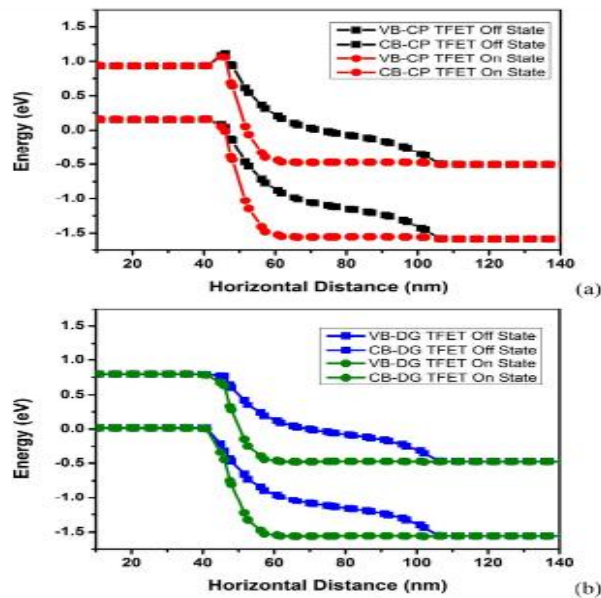


Figure 3. Variation of CV and value VB ranges for cavity lengths of 5 & 7 nm (a) $K = 2$ & $K = 10$ (b) $N_f = -5e11$ & $5e11$ with $V_{Gate-Source} = 0.5$ Volt, $V_{Drain-Source} = 0$ Volt.

Moreover, an increase in dielectric constant results in heightened band bending, leading to a reduction in barrier width (Narang et al., 2012). With cavity lengths set at 5 nm and 7 nm, Figure 3(c) illustrates alterations in the conduction and valence bands as cavity length varies, while increasing electron tunneling probability. Figures 2(b) and 3(b) demonstrate the significant influence of charged living particles on the device's bands. Biomolecules attached

to the cavity affect the barrier width between the source's valence band and the channel's conductive band(Chandan et al., 2018).

A result is obtained as an upsurge in opening length from 7 nano to 10 nm is approximately 0.0%.

3.2 Influence of uncharged and charged living biomolecules

For the variation with biomolecules present within the nanogap, leads to an increase in air's dielectric constant ($K = 1$).

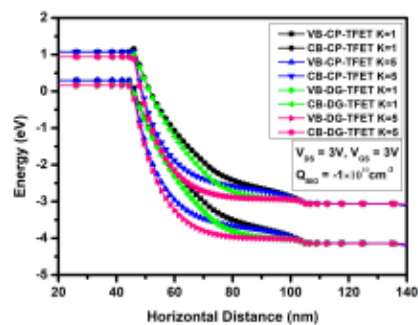


Figure 4. Surface potential variation with neutral biomolecules at a cavity length of 5, $V_{gs}=3$ V.

The filling of the cavity with biomolecules results in a significant rise in gate capacitance, rendering it an optimal sensor choice for CPB(Sant et al., 2018). However, as the dielectric constant increases, the effective capacitance decreases, impacting barrier width. Notably, immobilizing charged biomolecules in the gate underlap area induces remarkable fluctuations in surface potential(Mehrotra, 2016). For instance, when $K=6$ and $n_f=-110^{10}$ to 110^{12} cm^{-2} , the presence of positively charged biomolecules leads to increased surface potential. Moreover, the entry of neutral biomolecules into the cavity is facilitated.

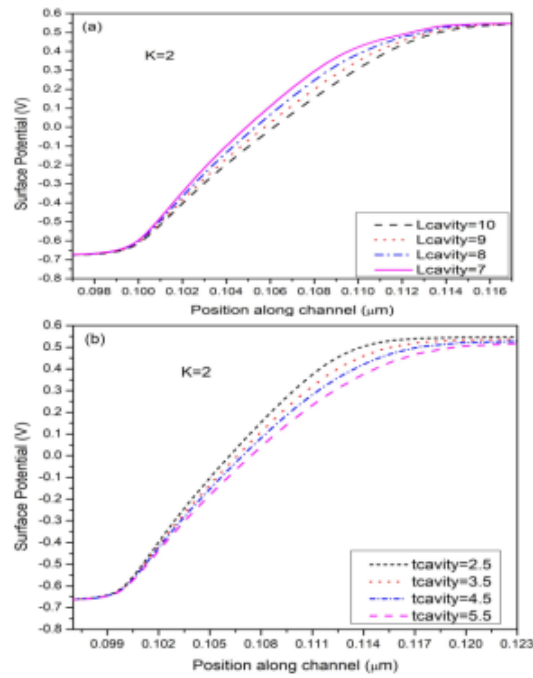


Figure 5. Surface potential variation at $V_{ds}=0$ V & $V_{gs}=0.5$ V.

Biomolecules with negative charges induce an elevation in flat band voltage within the overlap area, resulting in a reduction in effective gate bias (Samuel & Rao, 2022). Consequently, a decrease in surface potential occurs. Conversely, living elements with positive charge densities elicit a distinct response, as demonstrated in Figure 5(b).

Table 1.k of Different particles

Particles	K	Publication
Lysine	3.3	Yoon et al, (2020).(Yoon et al., 2020)
Arginine	3.2	Wadhwa & Raj, (2019).(Wadhwa & Raj, 2019)
Uricase	1.54	Wadhwa& Raj, (2018).(Wadhwa & Raj, 2018)
Protein	2.50	Intekhab Amin et al., (2016).(Anand et al., 2016)

The table presents dielectric constant values for select biomolecules, each with varying conductivity. Biomolecules interacting within the gate underlap area remain unutilized, with the gate underlap configuration for DM-JLFET under consideration.

3.3 Variation of Surface Potential with Cavity Length and Thickness

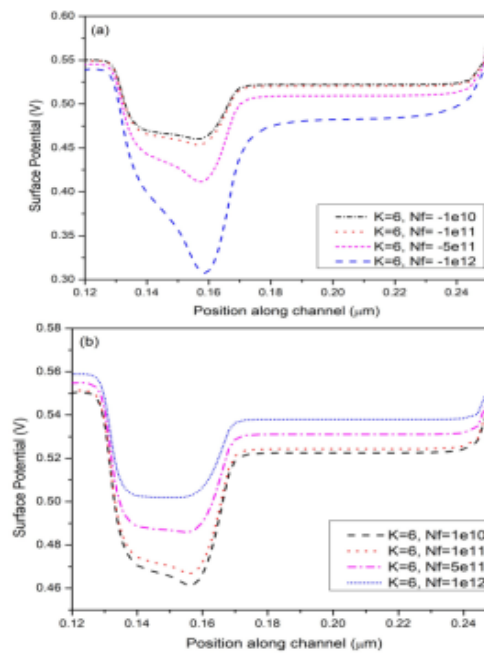


Figure 6. (a) different cavity lengths when $t_{cavity} = 2.5 \text{ nm}$ (b) $L_{cavity} = 5 \text{ nm}$.

The widening of the barrier due to cavity formation results in a decrease in surface potential. As cavity length increases, the device's surface potential diminishes, with a 7-nm-long cavity exhibiting lower surface potential than a 10-nm-long one (Ajay et al., 2013). Moreover, an increase in gate overlap area thickness leads to a drop in surface potential. As the opening escalates, the surface potential variance reduces. Notably, the maximum surface potential is achieved when the cavity is 2.5 nm thick. Fluctuations in surface potential occur at the source-channel junction due to biomolecule conjugation, particularly with a channel length set at 7 nm.

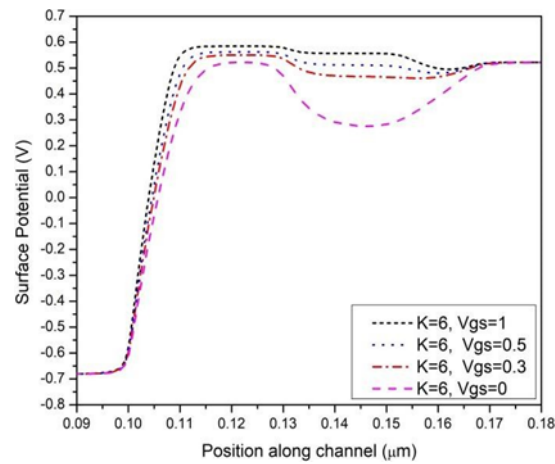


Figure 7. Surface potential variation sensitivity characteristics.

As gate bias increases, the gate-to-channel coupling weakens. However, with continued rise in gate bias, the two behaviors converge, indicating complete gate control. Approximately one million volts in surface potential change corresponds to one million volts in gate bias voltage. Comparative analyses of surface potential investigations highlight significant alterations due to biomolecule sample characteristics under suitable biasing conditions.

3.4 Surface Potential Sensitivity Variance

Biosensors' efficacy can be evaluated through their surface potential sensitivity, gauging their responsiveness to positive and negative potentials. Surface potential sensitivity is calculated using the formula: $(\text{bio}-z)/z$, which equals $(100\% \text{ bio}-b)$ where z represents volumetric potential. When the cavity is vacant, surface potential values for z and z_{bio} are: $z = 0$ and $z_{\text{bio}} = 0$.

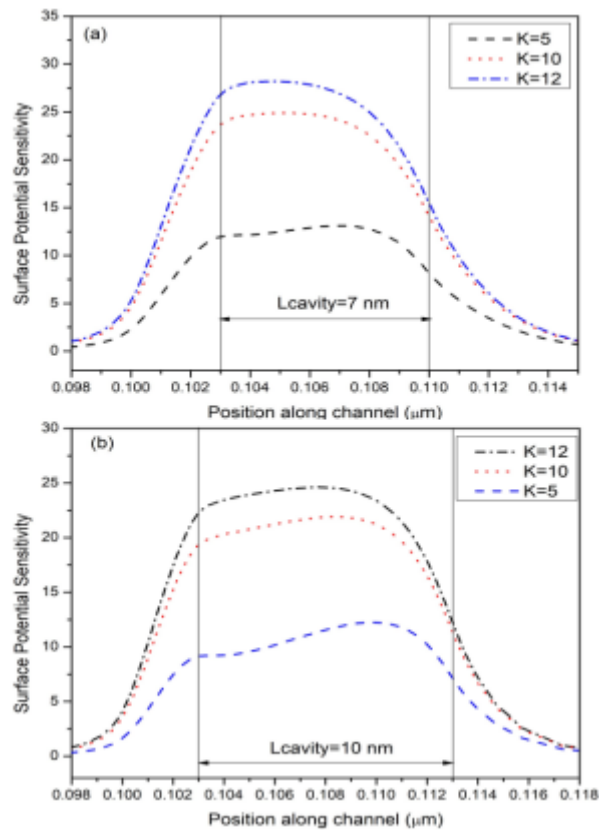


Figure 8. (a) The opening is 7 nm and (b) 10 nm.

In Figure 8, parameters are presented as a function of cavity length, utilizing a dielectric constant with V_{gs} and V_{ds} equal to 0.3v and 0v, respectively (Park et al., 2012). When biomolecule concentration increases, biomolecules immobilized within the cavity enhance surface potential sensitivity concerning dielectric constant change. The CPB DM-JLTFET device's biomolecule detection capability within the cavity correlates.

The incidence of +ve living particles lessens barrier thickness, consequently enhancing surface potential sensitivity. Conversely, the effect of charged living particles on surface potential compassion substantially upsurges with proteins that are +ve, while it diminishes with negatively charged biomolecules in the gate overlap zone. Increased charged biomolecules result in barrier band bending, further decreasing barrier width and significantly augmenting surface potential sensitivity.

3.5 Sensitivity of Ion/Ioff Ratio to Cavity Length Fluctuation

The impact varies with the opening. Ion state current begins to decrease as gate capacitance diminishes.

Table 2a. Cavity Length Change at Vgs and Vds

Cavity Length	Vgs=1.5v, Vds=0.5v	Vgs=1.5v, Vds=1.5v
	Ioff	Ion
5 nm	3.89E-18	6.08E-11
7 nm	3.91E-18	6.31E-13

Table 2b. Cavity Length Change at Vgs and Vds

Cavity Span	1.5 VoltV_{Gate-Source}, 0.5 VoltV_{Drain-Source}	Vgs=1.5, Vds=1.5
	Ioff	Ion
5 nm	3.85E-18	1.68E-10
7 nm	4.32E-18	3.19E-12

The results presented in Table III indicate the influence of neutral biomolecules, termed ion/Ioff ratio. K=1, 2, 3, 6, 10 exert different effects on the Ion/Ioff ratio of the devices. An increase in the Ion/Ioff ratio occurs as barrier width decreases. Ion/Ioff ratio serves as a determinant of the sensor's sensitivity. However, the Ion/Ioff ratio varies once the opening is 5 nm and 7 nano for the CPB DM-JLTFET device.

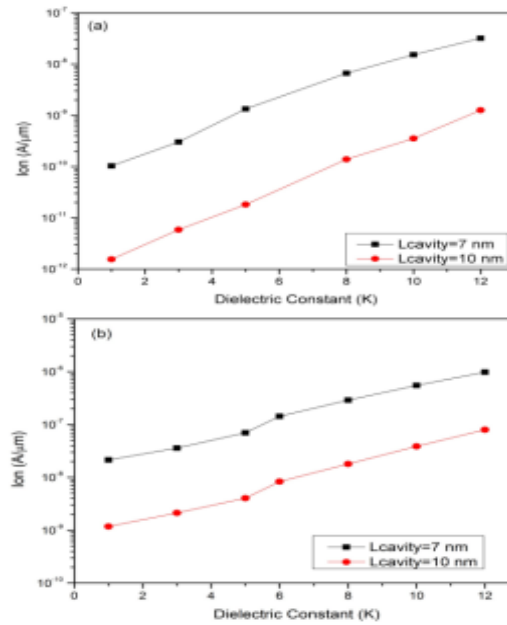


Figure 9. Plots of Ion when (a) 1.5 Volt $V_{Gate-Source}$ and 0.5 Volt $V_{Drain-Source}$ (b) $V_{Gate-Source}$ and $V_{Drain-Source}$ 1.5 Volt.

3.6 Sub-threshold slope (SS)

Another crucial parameter for examination is the sub-threshold slope (SS) Unlike thermionic emission, the SS relies on tunneling principles ($V_{ds} = 0.5v$ and $1.5v$). As the dielectric constant increases, the value of SS begins to decrease. This indicates that the SS parameter reaches a minimum at 7 nm, compared to 10 nm. Fig. 10 illustrates the decrease in SS due to an increase in the dielectric constant.

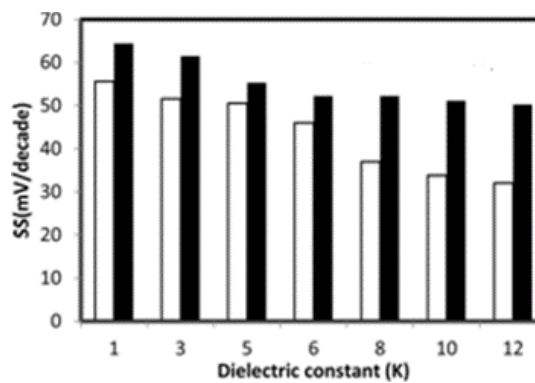


Figure 10. Plotting of SS when (a) V_{gs} 1.5V & V_{ds} 0.5V (b) at V_{gs} & V_{ds} 1.5V.

4. SUMMARY

The CPB DM-JLTFET structure dowries a highly effective approach for detecting immobilized biomolecules. With features like label-free detection, significant improvements in sensitivity, and a larger Ion/Ioff ratio, this device offers a cost-cutting clarification for the growth of biomedical diagnostic apparatuses. Through careful choice of cavity dimensionsadjoining the intersection of tunnelling under the right voltagesettings, maximum sensitivity is achieved. The evaluation of biomolecule effects on the electrical properties of the device reveals significantly lower leakage currents and enhanced sensitivity to charged biomolecules compared to existing FET devices, highlighting the superior performance of the CPB DM-JLTFET structure.

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